# REPORT DOCUMENTATION PAGE

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#### 14. ABSTRACT

We report the development of light emitters based on hyperbolic metamaterials. During the 18 month program period in Queens College of CUNY (Nov 2012 – May 2014), we successfully demonstrated growth of ultrasmooth silver films using germanium wetting layer, use of a high refractive index contrast grating to out-couple light from active hyperbolic metamaterials. We also successfully demonstrated for the first time simultaneous enhancement in spontaneous emission ad light extraction from active metamaterial structures.

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# **Report Title**

Final Report: Enhanced Light Emitters based on Metamaterials

#### **ABSTRACT**

We report the development of light emitters based on hyperbolic metamaterials. During the 18 month program period in Queens College of CUNY (Nov 2012 – May 2014), we successfully demonstrated growth of ultrasmooth silver films using germanium wetting layer, use of a high refractive index contrast grating to out-couple light from active hyperbolic metamaterials. We also successfully demonstrated for the first time simultaneous enhancement in spontaneous emission ad light extraction from active metamaterial structures.

Enter List of papers submitted or published that acknowledge ARO support from the start of the project to the date of this printing. List the papers, including journal references, in the following categories:

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03/29/2015 10.00	Xiaoze Liu, Tal Galfsky, Zheng Sun, Fengnian Xia, Erh-chen Lin, Yi-Hsien Lee, Stéphane Kéna-Cohen, Vinod M. Menon. Strong light–matter coupling in two-dimensional atomic crystals, Nature Photonics, (12 2014): 0. doi: 10.1038/nphoton.2014.304			
03/29/2015 9.00	T. Galfsky, H. N. S. Krishnamoorthy, W. Newman, E. E. Narimanov, Z. Jacob, V. M. Menon. Active hyperbolic metamaterials: enhanced spontaneous emission and light extraction, Optica, (01 2015): 0. doi: 10.1364/OPTICA.2.000062			
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03/29/2015	2.00	Vinod M. Menon, Harish Krishnamoorthy, Zubin Jacob, Evgenii Narimanov, Ilona Kretzschmar. Broadband QED using Hyperbolic Metamaterials, Conference on Coherence and Quantum Optics. 12-AUG-13, Rochester, New York.:,
03/29/2015	8.00	T. Galfsky, H. Krishnamoorthy, V.M. Menon, W. Newman, Z. Jacob, E. Narimanov. Extracting Light from High-K Modes in a Hyperbolic Metamaterial, 2014 IEEE Photonics Society Summer Topical Meeting Series. 14-JUL-14, Montreal, QC, Canada.:,
03/29/2015	6.00	Xiaoze Liu, Tal Galfsky, Fengnian Xia, Erh-chen Lin, Yi-Hsien Lee, Ashwin Ramasubramaniam, Stephane Kena-Cohen, Vinod M. Menon. Strong light-matter coupling in atomic monolayers, CLEO: QELS_Fundamental Science. 08-JUN-14, San Jose, California.:,
03/29/2015	5.00	Harish Krishnamoorthy, Ward D. Newman, Evgenii Narimanov, Zubin Jacob, Vinod M. Menon, Tal Galfsky. Directional emission from quantum dots in a hyperbolic metamaterial, CLEO: QELS_Fundamental Science. 08-JUN-14, San Jose, California.:,
03/29/2015	4.00	H. N. S. Krishnamoorthy, Z. Jacob, T. Galfsky, E. Narimanov, I. Kretzschmar, V. M. Menon. Optical topological transition in metamaterials: QED and related effects, 2013 IEEE Photonics Conference (IPC). 08-SEP-13, Bellevue, WA, USA.:,
08/30/2013	1.00	Zubin Jacob, Tal Galfsky, Evgenii E. Narimanov, Ilona Kretzschmar, Vinod M. Menon, Harish Krishnamoorthy. Topological Transitions in Metamaterials: QED and Related Effects, CLEO: QELS_Fundamental Science. 10-JUN-13, CLEO: QELS_Fundamental Science (2013).:,
TOTAL:		6

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**Sub Contractors (DD882)** 

**Inventions (DD882)** 

# **Scientific Progress**

The main scientific progress made during the program include:

- Realization of ultrasmooth sub-wavelength thick silver films for hyperbolic metamaterials
- Using high index contrast germanium gratings for light extraction
- Simultaneous enhancement in spontaneous emission rate and light extraction from active hyperbolic metamaterials.

Further details about the scientific accomplishments can be found in he attached document.

**Technology Transfer** 

Report Type: Final Report

Proposal Number: 62509EL

Agreement Number: W911NF1310001

Proposal Title: Enhanced Light Emitters based on

**Metamaterials** 

Report Period Begin

Date:

11/15/2012

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Date: 05/31/2014

### ABSTRACT

We report the development of light emitters based on hyperbolic metamaterials. During the 18 month program period in Queens College of CUNY (Nov 2012 – May 2014), we successfully demonstrated growth of ultrasmooth silver films using germanium wetting layer, use of a high refractive index contrast grating to out-couple light from active hyperbolic metamaterials. We also successfully demonstrated for the first time simultaneous enhancement in spontaneous emission ad light extraction from active metamaterial structures.

#### **OBJECTIVE**

To utilize hyperbolic metamaterials (HMM) to enhance the performance of LEDs and realize sub-wavelength lasers. Specific goals include:

- Develop ultrafast and bright LEDs using hyperbolic metamaterials.
- Demonstrate enhancement of light emission from inherently low quantum efficiency emitters such as silicon nanocrystals.
- Develop sub-wavelength size lasers using emitters embedded in hyperbolic metamaterials.

#### 1. Introduction

Under the proposed research program we are developing active HMMs comprising of one-dimensional metal-dielectric structures with colloidal quantum dots. Schematic drawing of the structure is

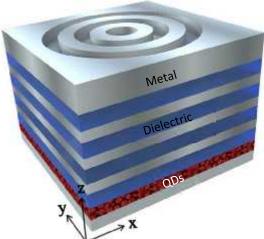


Fig. 1 Schematic of the metal dielectric metamaterial structure that results in the hyperbolic dispersion

shown in **Fig. 1**. The colloidal QDs are deposited via spin coating and capped with poly methyl methacrylate. The metal and dielectric layers are deposited via sputtering/electron beam evaporation. The gratings are defined via focused ion beam etching.

In any medium, the optical iso-frequency curve  $\check{S}(\vec{k}) = const$  can be engineered by tailoring the dielectric tensor  $\ddot{V}(\vec{r})$ . Metal dielectric composites can make the permittivity anisotropic and can considerably distort the topology of the iso-frequency curve. We consider the case of metal-dielectric composite metamaterials such as the one shown in **Fig. 1** which have a uniaxial form of the dielectric tensor  $\ddot{\delta} = diag(\dot{\delta}_{xx}, \dot{\delta}_{xx}, \dot{\delta}_{z})$  The iso-frequency curve for the extraordinary (TM-polarized) waves propagating in such strongly anisotropic metamaterial is given by:

$$\frac{k_x^2 + k_y^2}{\grave{o}_{\perp}} + \frac{k_z^2}{\grave{o}_{\parallel}} = \frac{\breve{S}^2}{c^2}$$

Since the length scale of the substructures is much smaller than the wavelength of light, one can define effective dielectric constants that control the macroscopic electromagnetic properties. The effective dielectric constants in the parallel and perpendicular directions for a one-dimensional metal-dielectric stack can be written as:

 $\frac{1}{\epsilon_{\parallel}} = \frac{f_a}{\epsilon_a} + \frac{f_b}{\epsilon_b}$  and  $\epsilon_{\perp} = f_a \epsilon_a + f_b \epsilon_b$ , where  $f_a$  and  $f_b$  are the fill fractions and  $\epsilon_a$  and  $\epsilon_b$  are the dielectric constants of the two materials, respectively. By appropriate choice of fill fractions and

the metallic and dielectric components, one has tremendous degree of control over the effective dielectric constants of the metamaterial structure. Closed iso-frequency surfaces different from a simple sphere (eg: ellipsoid) can occur in these metamaterials when  $\epsilon_{\parallel} > 0$  and  $\epsilon_{\perp} > 0$ . On the other hand an extreme case of iso-frequency surface modification occurs when the dielectric constants show opposite sign ( $\dot{q} < 0$  and  $\dot{q}_{\perp} > 0$ ) such that the iso-frequency curve opens up into a hyperbolic surface. This can be accomplished by controlling the fill fraction of the metal. The effective dielectric constants of one such metal dielectric structure as a

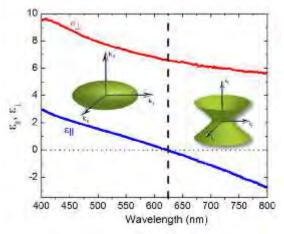


Fig. 2 Dispersion of a strongly anisotropic metamaterial exhibiting the elliptical and hyperbolc dispersion regimes.

function of wavelength is shown in **Fig.2** The real part of the effective dielectric constant parallel to the layers (blue) goes through a sign change at 620 nm, while the perpendicular component (red) stays positive. This results in the hyperbolic dispersion. The shape of the dispersion curves are shown in the inset of **Fig. 2** for the different spectral regions.

One of the major issues in realizing HMMs with desired dispersion is the control over the layer thickness and the smoothness of the metal films. Unfortunately, at thickness less that 20 nm, silver films are usually percolated leading to island formation. This in turns affects the propagation length of surface plasmons at the metal-dielectric interface. Hence, one of the first issue we addressed under the program was to develop techniques to realize ultra-smooth silver films.

#### 2. Project Descriptions

# Growth of ultra-smooth silver films for layered hyperbolic metamaterials:

Growing thin (<20 nm) silver films on dielectric surfaces result in island formation since the thickness is below the percolation threshold. Scanning electron microscope image of a typical silver film, 10 nm in thickness grown on a glass substrate is shown in Fig. 3a. This film is highly percolated and hence mostly support localized plasmon modes. More recently, following and tweaking approaches reported previously [1, 2], we have been successful in demonstrating growth of ultra-smooth silver films on dielectric surfaces. The key here is to use a wetting layer which in our case was Ge. The structure consisted of Glass substrate/ Al<sub>2</sub>O<sub>3</sub>/Ge/Ag. Scanning electron microscope image of the silver film grown using the Ge wetting layer is shown in Fig. 3b clearly indicating a smooth surface. Fig. 3c shows the atomic force microscope image of the surface of the. Less than 200 pm r.m.s roughness is observed in the Z (growth) direction and less than 2 nm roughness the plane. r.m.s in х-у

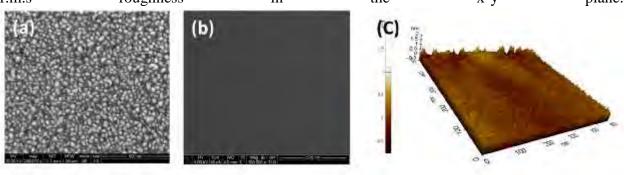


Fig. 3 Scanning electron microscope image of surface of slver film grwn (a) without a wetting layer and (b) with a wetting layer of Ge. (c) Atomic Force Microscope image showing ultra-smooth surface topography of the film grown using the wetting layer.

#### Design and fabrication of bulls eye grating structures for better in/out coupling:

One of the limiting aspects of HMM structures is the difficulty in coupling light into and out of the structures, thus making them difficult to implement in practical photonic device configurations. This issue arises due to the presence of high-k stats in HMM structures which cannot couple to light in vacuum. One way to circumvent this issue while preserving the high-k states is to use a grating structure that couples light of specific k-vectors into and out of the structure. In this context we have recently designed a bulls eye grating etched on to the metallic layer which efficiently out-couples the high-k states from the HMM into free space. Schematic drawing of the grating structure is shown in **Fig. 1**. Design of the grating structures were carried out using COMSOL finite element modeling package. Show in **Fig. 4** is the simulated electromagnetic field profile along with the scattering when a radiating dipole is placed in the near field of the HMM. In the absence of the grating structure most of the light either gets absorbed or gets reflected (**Fig. 4a**). On the other hand when a grating is placed on the top layer, one can convert the in-plane surface

plasmon modes that are evanescent in nature into propagating modes (**Fig. 4b**). Shown in **Fig. 5** (**a,b**) is the optimization of the grating period to out-couple a specific mode out of the HMM structure. Scanning electron microscope image of a bulls eye grating fabricated using focused ion beam etching is shown in **Fig. 5c**. The etching was carried out on structures where the silver film was not optimized for smoothness and hence one can see the rough surface on the grating. In the next phase of the project we are fabricating gratings on the optimized smooth silver films.

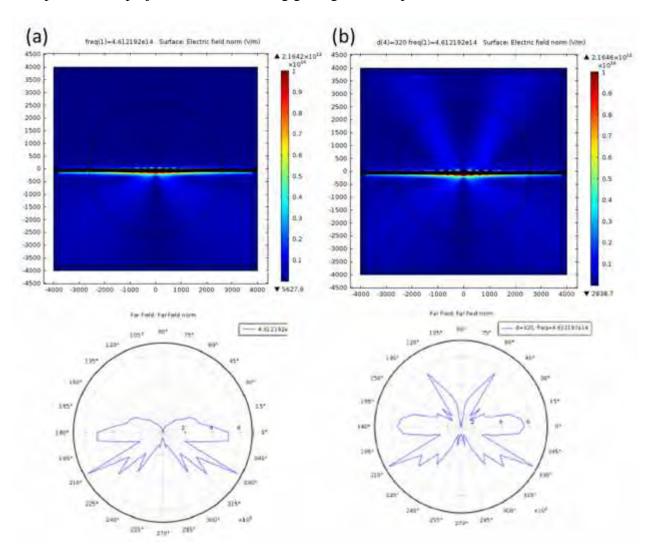
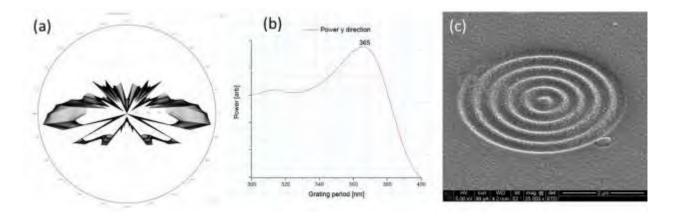


Fig. 4 Simulated electric field profile and scattering directions for a dipole placed in the near field of a HMM structure. (a) with no grating and (b) with a grating. It is clearly seen that in the presence of a grating, there is forward propagating waves that can be detected in the far field.



**Fig. 5** Optimization of the grating period for maximum forward scattering amplitude. (a) polar plot showing the scattering angles and (b) the power at far field as a function of grating period. In this specific case for wavelength of 650 nm, we obtain maximum efficiency in outcoupling with a grating of 365 nm period. (c) SEM image of grating fabricated on a HMM structure using focused ion beam technique.

# Fabrication of HMMs with ultrasmooth silver films and out-couplers

Using the results from the above two projects we fabricated the **HMM** structures with smooth silver films and germanium out-couplers. Shown in Fig. 6a is the schematic of the strcture along with the cross-sectional TEM images of the layers showing smooth silver and slumina layers. The quantum dots are emdeeded inside the HMMs, Shown in Fig. 1c is the effective dielectric constants of the HMM strcture realized. In the vicinity of the CdSe quantum dot emisison, metamaterial cleraly has one component of its dielectric constant negative (ε||) and hence results in hyperbolic dispersion. Shown in Fig. 1d is the measured spontaneous emision lifetime of quantum dots on glass, a control sample of 1 unit cell, a 4 period strcture and a 7 period strcture where the quantum dots are embedded inside. Clearly the 7 period strcture shows the most reduced lifetime as well as the naoowest distribution in lifetime. The reduced lifetime is due to enhancement in spontaneous emission rate. The reduced distribution in lifetime is

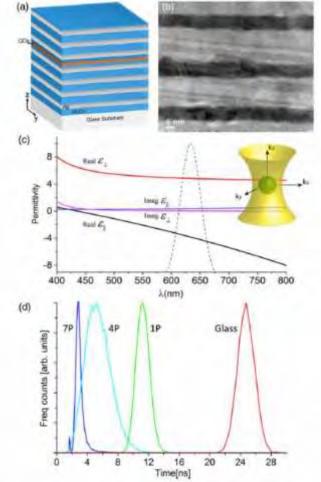
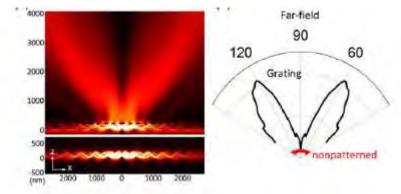


Fig. 6 (a) Schematic of the active HMM structure and (b) the cross sectional TEM image of the structure showing the alumina (light) and silver (dark) layers. (c) Effective dielectric constants of the active HMM structure. (d) Lifetime distributions of CdSe quantum dots embedded in different structures.

because the quantum dots being embedded inside the HMM allow almost identical coupling to the high-k modes of the HMM.

Finite element simulations using COMOSL were caried out to design the out-coupler. Shown in **Fig. 7** is the simulation showing the outcoupling from the optimzed grating. Also for comparison the emission in the absence of a grating is also shown.

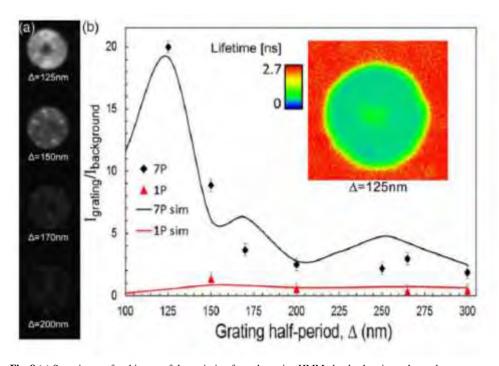


Finally in **Fig. 8** (a) we show the emission from active HMMs with

Fig. 7 (a) results of finite element simulations showing outcoupling using the bulls-eye grating structure.

different pitch gratings on top captured via confocal imaging. We see the 125 nm pitch to yield the most light outut. The samples weer pumped with a 440 nm laser and imaged using a scanning

confocal imaging set up with an avalanche photo diode (APD) detector. **Fig. 8(b)** shows the ratio of the emission intensity from grating the region to the background (unpatterned) of the active HMM. Clearly we see a factor of 20 enhancement emision intensity for the 125 nm itch grating. Shown in the inset of Fig. 8(b) is the lifetime map of the emision taken using a fuoroscence lifetime imaging set



**Fig. 8** (a) Scanning confocal image of the emission from the active HMM clearly showing enhanced emission from the 125 nm pitch grating structure. (b) Ratio of measured emission intensity at the grating to that of the background as a function of grating period. Inset: Fluorescence lifetime image of the structure showing decreased spontaneous emission lifetime from the grating region.

up. The lifetime of the emission is found to be shortest within the grating region. This is due to the more efficient coupling of vertical dipole emission which undergoes the greater Purcell enhancement to the grating. This work was published in *Optica* (2015) [3].

## Design of sub-wavelength cavities using HMMs

In this part of the project we used finite difference time domain simulations to design subwavelength sized cavities made of HMMs that can efficiently confine light at high-k values. A typical structure in this case is shown in **Fig. 9a**. The structure consists of alternating layers of Ag and TiO<sub>2</sub>. The total thickness of the structure is 30 nm and the lateral dimensions are 45 nm x 45 nm. This structure is shown to support a resonant

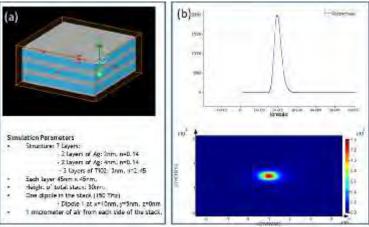


Fig. 9 (a) Schematic of the simulated sub-wavelength cavity structure along with simulation parameters and (b) simulation results showing resonant mode at  $2\mu m$  and the associated intensity profile inside the cavity at the resonance.

mode at a wavelength of 2  $\mu$ m (**Fig. 9b**) which is far larger than the physical dimensions of the structure. The electric field inside the HMM cavity along the different directions is shown in **Fig.** 10. It is clearly seen than the mode is well confined within the sub-wavelength HMM structure along all three dimensions.

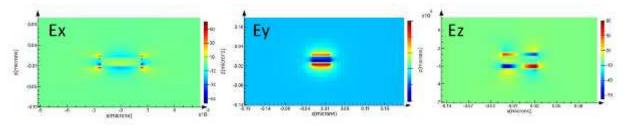


Fig. 10 Electric field simulations showing well confined modes in all three directions at 2µm in the sub-wavelength structure.

#### 3. REFERENCES

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